

## 2D MATERIALS

## Brightening the dark excitons

Two independent methods using near-field coupling to surface plasmon polaritons and magnetic brightening allow the observation of dark excitons in  $\text{WSe}_2$ .

Kian Ping Loh

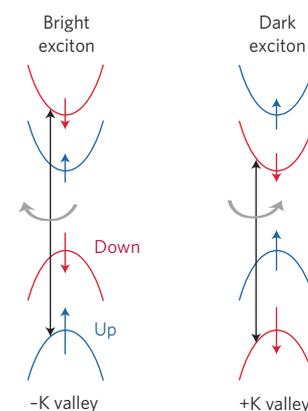
Optical properties of semiconductors are dominated by excitons, which are electron–hole pairs bound by Coulomb interactions. Excitons can be bright or dark depending on the spin orientation of the individual carriers. In a bright exciton the electron and the hole have antiparallel spins and can recombine easily through the emission of a photon. In a dark exciton the spins are parallel and the recombination cannot occur via direct emission of a photon as this would not allow for spin momentum conservation. Thus, dark excitons have much longer radiative lifetimes than bright excitons. This large difference in recombination time can potentially be harnessed to create ‘fast’ or ‘slow’ light. Besides, the highly stable, non-radiative nature of dark excitons makes them attractive for optically controlled information processing.

To take advantage of the features of dark excitons, it is essential to create conditions in which these excitonic species can exist. It turns out that the formation of dark excitons occurs in  $\text{WS}_2$  or  $\text{WSe}_2$ , where rotational and inversion symmetry-breaking together with strong spin–orbit coupling leads to spin-split energy states. In these materials dark excitons are lower in energy than the bright excitons in the conduction band<sup>1</sup>. However, in order to access the stored spin information, the dark exciton needs to have the ability to respond to light. Making use of the optical selection rules in 2D transition metal dichalcogenides (TMDs)<sup>2</sup>, it is possible to selectively control spin and valley excitation to initialize this state (Fig. 1). This has piqued the interest of researchers and efforts are ongoing to find ways to optically detect and manipulate dark excitons.

Writing in *Nature Nanotechnology*, two teams, Xiao-Xiao Zhang and co-workers<sup>3</sup>, and You Zhou and co-workers<sup>4</sup>, independently describe two approaches to induce light emission from dark excitons in monolayer  $\text{WSe}_2$ , thus opening new possibilities to store and manipulate valley and spin information.

Applying a magnetic field to a semiconductor can mix electronic wave functions and shift the spectral weight between bright and dark excitons depending on the magnetic field direction<sup>5</sup>. An in-plane magnetic field does not couple to the in-plane motion of charge carriers in 2D materials and hence, does not perturb the electronic structure. Instead, it alters the spin alignment of carriers, which relaxes the spin-selection rule and makes dark excitons optically detectable by photoluminescence (PL) spectroscopy. Zhang *et al.*<sup>3</sup> demonstrated that by applying an in-plane magnetic field greater than 10 T, a magnetic brightening of the dark exciton can be achieved in  $\text{WSe}_2$ . Using circularly polarized light to excite the dark state, the researchers show that the dark exciton is defined by its spin state as well as the chiral information characteristic of the valley it resides in. By measuring the polarization of the light emitted from the dark exciton, the team discovered that the dark exciton, in addition to having an opposite chirality to that of the bright exciton, also has a much longer valley recombination lifetime. Although the exact mechanism is not clear, Zhang *et al.* suggest that it could be due to the transfer of electrons from one photo-excited valley to another. This result in particular could have implications for valley-selective spin injection. Besides, because dark excitons in  $\text{WSe}_2$  emit light of certain chirality, this process is also attractive for developing photon emitters for chiral optics<sup>6</sup>. It is worth investigating if, instead of applying a high magnetic field, a similar effect can be achieved by coupling the 2D TMD to a ferromagnetic layer.

Although magnetic brightening is fundamentally important for the understanding of, for example, the structure of the spin-split conduction bands in 2D TMDs, it is not easy to make use of this effect in conventional device fabrication schemes. Luckily, an alternative way of probing the dark exciton is available, which simply involves placing  $\text{WSe}_2$  monolayers on patterned silver



**Figure 1** | Schematic illustration of bright and dark excitonic states. Bright excitons (left) are optically allowed and dark excitons (right) are optically forbidden. Bright and dark excitons can have different chirality due to valley separation. The black arrows indicate photons with different chirality. Blue and red lines indicate levels containing ‘up’ and ‘down’ spins, respectively.

surfaces — the approach pioneered by Zhou and colleagues<sup>4</sup>. Silver surfaces have a high density of free electrons, which can be collectively excited to form plasmons. When light interacts with these species, a combined electromagnetic radiation called the surface plasmon polariton (SPP) is generated. Coupling with SPPs can be used to selectively enhance optical transitions with dipole moments normal to the 2D plane. Dark excitons in  $\text{WSe}_2$  have a nonzero dipole moment along the  $z$  direction, allowing them to couple to the SPP modes, which are also  $z$ -polarized. This ‘switches on’ the PL response from the dark excitons, yielding a PL intensity equivalent to that achieved by a 30 T in-plane magnetic field.

The advantage of the SPP approach is its ease of integration with atomically thin 2D materials. Silver mirrors can be patterned in the form of subwavelength structures called metasurfaces to create optical interference effects, which enables spatial variation of the optical response

from  $\text{WSe}_2$  coupled to the silver mirror. It is worth noting that since such excitonic effects enhance optical nonlinear properties by orders of magnitude, regions with locally varying nonlinear effects can be fabricated by using metasurfaces to control the brightening of dark exciton states on a 2D monolayer (for example,  $\text{WSe}_2$ ). Going forward, optical schemes based on quantum coherence<sup>7</sup> and multiphoton excitation<sup>8,9</sup> can be used to manipulate the dark excitons on such metasurfaces to

generate light with different frequencies, enabling quantum information processing and storage.

In future nanophotonic devices, 2D materials can be integrated into a photonics platform to make optical circuitry in which dark excitons can be used for encoding and transporting information on a chip.  $\square$

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## NANOLASING

# Multimode superlattice arrays

A superlattice of gold nanoparticles shows stable lasing at different frequencies.

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**M**ultifrequency lasers are currently used in communications, sensing and measurement technologies as they provide multiple channels and enhanced power. But optical devices of this kind are challenging to produce at the nanoscale, because multiple modes usually compete over the available gain, making them unstable. Writing in *Nature Nanotechnology*, Wang *et al.* now present a device that shows multimode lasing at the nanoscale<sup>1</sup>. Lasing comes from the emission from modes at the energy band edges of a superlattice array of gold nanoparticles. Because the superlattice can be designed to sustain several such modes, the device can allow lasing at different frequencies.

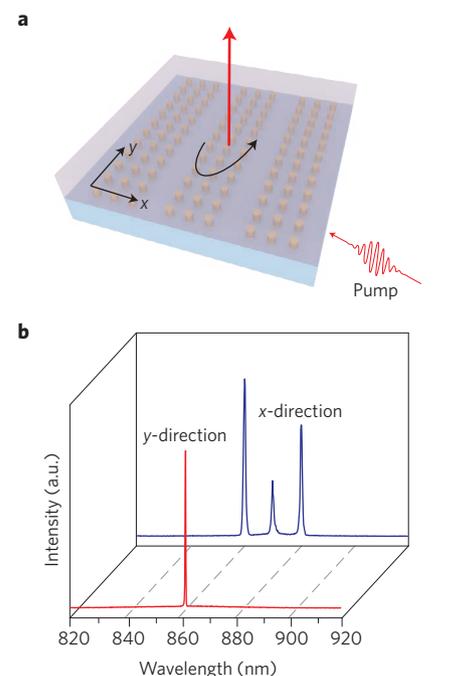
Nanostructured metal films as well as nanoparticles can sustain hybrid modes of surface-bound electromagnetic fields and electron plasma oscillations<sup>2–4</sup>. Noble metal nanoparticles, such as gold and silver, possess prominent plasmonic resonances that make them act like tiny antennas, radiating light and creating a strong electromagnetic field (or hot spot) around them. It has been suggested that these resonances can show lasing<sup>5</sup>, and indeed a few examples of plasmonic lasers exist<sup>6</sup>. In particular, plasmonic band-edge lasing in periodic arrays of metal nanoparticles<sup>7–9</sup> offers narrower linewidths than typical plasmonic lasers, as well as the freedom to design the band structure at will by changing the geometric parameters of the array. Wang *et al.* extend this concept to plasmonic superlattices.

The researchers arrange gold nanoparticles (120 nm in diameter) on

a glass substrate in a periodic array such that the interparticle distance is 600 nm in 18- $\mu\text{m}$ -long patches, with the patches separated from each other by 24  $\mu\text{m}$  (Fig. 1). This superlattice is then embedded within a solution containing an organic dye (IR-140-DMSO) as gain medium. The dye molecules are excited by pumping with 90-fs laser pulses. The emission from the sample shows several footprints of lasing, such as a line narrowing, both in frequency and momentum, and a threshold pump power for the onset of lasing. Remarkably, the researchers observe lasing at several frequencies simultaneously. The lasing frequencies and the number of lasing modes can be controlled by changing the lattice design and the excitation conditions.

An attractive feature of this system is that the multiple lasing action is very stable. According to simulations, this is because different modes in the superlattice array have field intensity maxima in distinct, non-overlapping regions of space. For instance, one mode could mainly be located closer to the nanoparticles whereas another could lie in the area between them. Therefore, these modes have no strong, direct competition for gain, because excited dye molecules in different parts of the sample primarily emit to different modes.

In addition, by fabricating an array with superlattice periodicities only in one direction, Wang *et al.* could switch between single and multiple lasing modes by simply rotating the sample with respect to the excitation direction. To take full advantage of this single-to-multiple mode switching, an obvious research direction



**Figure 1** | Multimode nanoscale lasing. **a**, Superlattice arrays of gold nanoparticles on glass embedded with a gain medium of organic molecules. **b**, The superlattice seen in the x-direction contains multiple band edges providing lasing at multiple frequencies simultaneously (blue curves). In contrast, a single-period array (y-axis) leads to one lasing peak (red curve). By rotating the sample (black arrow in **a**) with respect to the pump polarization direction, switching between single-mode and multimode lasing is observed. Figure reproduced from ref. 1, Macmillan Publishers Ltd.